SCIENTIFIC REPORTS

Received: 25 January 2017 Accepted: 19 April 2017 Published online: 26 May 2017

OPEN Adaptive stochastic resonance for unknown and variable input signals

Patrick Krauss^{1,2}, Claus Metzner², Achim Schilling^{1,2}, Christian Schütz¹, Konstantin Tziridis¹, Ben Fabry² & Holger Schulze¹

All sensors have a threshold, defined by the smallest signal amplitude that can be detected. The detection of sub-threshold signals, however, is possible by using the principle of stochastic resonance, where noise is added to the input signal so that it randomly exceeds the sensor threshold. The choice of an optimal noise level that maximizes the mutual information between sensor input and output, however, requires knowledge of the input signal, which is not available in most practical applications. Here we demonstrate that the autocorrelation of the sensor output alone is sufficient to find this optimal noise level. Furthermore, we demonstrate numerically and analytically the equivalence of the traditional mutual information approach and our autocorrelation approach for a range of model systems. We furthermore show how the level of added noise can be continuously adapted even to highly variable, unknown input signals via a feedback loop. Finally, we present evidence that adaptive stochastic resonance based on the autocorrelation of the sensor output may be a fundamental principle in neuronal systems.

Biological organisms, as well as technical systems, rely on sensors that transmit environmental signals into the system for subsequent information processing. Sensors, in general, have a limited sensitivity, so that input signals with amplitudes below a certain threshold cannot normally be detected. Stochastic resonance (SR), a phenomenon first described by Benzi et al. in 1981¹ enables non-linear systems to detect even sub-threshold signals by means of added noise²⁻⁴, which brings the weak input signal at random time points to above-threshold levels. This mechanism, however, requires a tuning of the added noise level to work properly. In the case of input signals with known properties, such as periodic signals with known frequency, the optimum noise level can be found by maximizing the signal-to-noise ratio of the sensor output^{4,5}. For arbitrary, non-periodic input signals, however, it is not possible to separate signal and noise in the detector output by frequency filtering, so that the quality of the signal transmission must be assessed in a different way.

Regardless of the detailed mechanism, SR has been identified in a wide range of biological systems, including the mating behavior of Nezara viridula that are able to detect subthreshold acoustic calling signals when mixed with acoustic Gaussian noise of suitable intensitiy⁶, rat cutaneous mechanoreceptors when stimulated with subthreshold aperiodic stimuli plus noise² and the paddlefish which relies on electrical signals, amplified with stochastic resonance, to hunt edible plankton⁷. SR has recently received increasing attention especially in the context of experimental and computational neuroscience where it helps to explain how neuronal systems operate in noisy environments⁸⁻¹⁰.

It can be demonstrated that an optimal level for the added noise exists that maximizes the information transmission from the sub-threshold input to the sensor output^{11,12}. In self-adaptive non-linear signal detection systems based on SR, the optimum noise level can be continuously adjusted (increased or reduced until an optimum is found) via a feed-back loop, so that the system response in terms of information throughput remains optimal even if the properties of the input signal change (Fig. 1). For this processing principle, the term adaptive SR has been coined¹³⁻¹⁵

A widely used measure (in the following refered to as 'objective function') for quantifying the information throughput is the mutual information (MI, cf. Methods) between the input signal and the sensor output^{3, 12, 14} a statistical quantity that measures the mutual dependence of these two variables¹⁶. It has been shown previously that the MI has a maximum at a well-defined, optimal intensity of the added noise¹².

To calculate the MI (or other objective functions such as the signal-to-noise ratio, or the cross-correlation between input and output signal), the input signal must be known^{2-4, 6, 10-13, 17}, but this is often not the case. Here,

¹Department of Otorhinolaryngology, University Erlangen, Nürnberg, Germany. ²Department of Physics, University Erlangen, Nürnberg, Germany. Correspondence and requests for materials should be addressed to H.S. (email: holger.schulze@uk-erlangen.de)



Figure 1. Principle of adaptive SR. Stochastic resonance describes how a sub-threshold signal with added noise of sufficient amplitude can result in a non-zero sensor output. If the input signal amplitude changes over time, the optimal noise level can be estimated and adjusted in a feedback loop from a resonance curve. If the probability density of the input signal levels is known, the resonance curve can be obtained from a variety of objective functions. For unknown input signals, however, the only applicable objective function is the autocorrelation of the sensor output.

we consider the practically important case of unknown and not necessarily periodic input signals. Although finding the optimal level of noise becomes less important with arrays of transducers¹⁷, it still remains an unsolved problem for single detector systems.

In this letter we show that this fundamental limitation of adaptive SR can be overcome by another objective function, namely the autocorrelation (AC, cf. Methods) of the detector output. Maximizing the output AC leads to similar or identical estimates of optimal noise intensities for SR as estimates based on established objective functions, yet with the decisive advantage that no knowledge of the input signal is required (Fig. 1).

Results

To demonstrate the equivalence of the mutual information (cf. Methods) between input and output, and the autocorrelation of output (cf. Methods), for finding the optimal noise level, we consider a so-called discrete-symmetric model with a symmetric threshold $|\theta| = 1.1$ and a discrete output $y_t = \{-1.0, 0.0, 1.0\}$, i.e. the sensor output is zero for a sensor input s_t of $-\theta < s_t < \theta$, -1.0 for $s_t < -\theta$, and 1.0 for $s_t > \theta$. As a discrete input signal we generate a correlated bipolar string $s_t = \{-1.0, 1.0\}$ with the probability of successive values being identical $p(s_t = s_{t+1}) = 0.7$ (Fig. 2 top). Without added noise, the input signal alone never passes the thresholds.

Furthermore, we introduce the concept of the success probability (cf. Supplements): i.e. the conditional probability that, given a certain input, the correct corresponding output is produced. For a discrete-symmetric sensor with correlated bipolar input, we prove analytically (cf. Supplements) that the success probability as a function of noise intensity has a well-defined peak, indicating the existence of an optimal level of noise for SR. We also prove that the mutual information and the autocorrelation of the output can be expressed as strictly monotonous functions of the success probability (cf. Supplements). Hence both, the mutual information and the autocorrelation of the output exhibit their maximum at the same level of noise. We numerically tested our approach and confirmed that both the mutual information and the autocorrelation of the output yielded similar optimal noise levels (Fig. 2 bottom). In addition, we applied our analytical model to the important case of soft thresholds¹⁸ and non-Gaussian noise and show that this only changes the success probability *Q*, whereas the MI and AC remain monotonous functions of *Q* and, hence, peak at the same level of noise (cf. Supplements).

To generalize our finding and demonstrate that both approaches for quantifying the information throughput give similar optimal noise intensities, we simulated a variety of models comprising different discrete and



Figure 2. Results of analytical model and numerical simulations. Top: Example of a discrete signal consisting of a correlated bipolar string with $p(s_t = s_{t+1}) = 0.7$. Bottom: Comparison of different objective functions versus noise intensity for correlated bipolar strings with symmetric threshold and a signal-to-threshold distance of 0.1. Analytical and numerical solutions coincide, and all objectives functions show a maximum at the same noise level. MI = mutual information, AC = autocorrelation of the output signal.

continuous input signals, different types of detector models and different signal-to-threshold distances (the difference between the temporal averaged sub-threshold signal and the threshold).

As detectors, we tested the above described discrete-symmetric model (Fig. 3d), a discrete-asymmetric (Fig. 3c), a continuous-asymmetric (Fig. 3a) and a continuous-symmetric (Fig. 3b) detector model (cf. Methods). These detectors are all memoryless, i.e. their output at a certain time depends only on their input at that time but neither on their internal state, nor on past inputs. We also tested a model with memory and implemented the biologically important neuronal leaky integrate-and-fire model (cf. Methods).

As discrete aperiodic input signals, we used the correlated bipolar string described above (Fig. 2 top). In addition, we tested different time series of continuous aperiodic signals, namely the time series x(t) of the 3-dimensional Roessler attractor (Fig. 4a top, cf. Methods), an Ornstein-Uhlenbeck process (Fig. 4b top, cf. Methods), and wave files of recorded speech (Fig 4c top). Furthermore, a sine wave with constant frequency and amplitude was used as continuous periodic input (Fig. 4d top).

Four sample resonance curves (i.e. the mutual information and the autocorrelation of the output as functions of noise intensity) are shown in Fig. 4. In all cases the mean signal-to-threshold distance in arbitrary units was set to 1.1. Although the resonance curves from the mutual information (red) and the autocorrelation of the output (blue) are different, they peak at identical noise intensities.

We then tested all combinations of the five different detector models (that is: continuous-symmetric, continuous-asymmetric, discrete-asymmetric, leaky integrate-and-fire) and the three input signal types (correlated bipolar string, Roessler attractor, sine wave) with twenty different signal-to-threshold distances (from 0.0 to 1.0 in steps of 0.05, arbitrary units) (Fig. 5). For each combination, ten trials were performed with different seeds for the random number generator resulting in a total of 3000 simulations. For each simulation, the optimal noise intensities according to the mutual information (x-axes) and the autocorrelation of the output (y-axes) have been evaluated and averaged over the ten performed trials per combination. Almost all tested combinations yielded nearly identical estimates of optimal noise intensity. The total correlation coefficient between optimal noise intensity based on mutual information versus autocorrelation was r = 0.96 with a mean squared error of e = 0.026. Except for the combination of a correlated bipolar string with the leaky integrate-and-fire neuron, all other combinations yielded correlation coefficients from 0.925 to 0.999 with mean squared errors from 0.010 to 0.042, further demonstrating the equivalence of the two measures.

Discussion

We have shown that the autocorrelation of the output signal can serve as a universal objective function to estimate the optimal level of noise in SR-based sensory systems. This new approach allows for the first time the technical



Figure 3. Detector models. Sketch, demonstrating the reponse of different sensor types to a sinusoidal input siganl (red line). We compare the four cases of continuous (**a**,**b**) or discrete (**c**,**d**) output (blue lines or dots, respectively) and with asymmetric (**a**,**c**) or symmetric (**b**,**d**) thresholds (green lines).

implementation of adaptive SR in situations where the information content of the input signal is unknown or highly variable. Once the signal is optimally transmitted by the sensor, the (noisy) output signal can be further improved by several techniques, such as kernel density estimation¹⁹.

Several interesting phenomena are closely related to stochastic resonance. A well-known technical system to reduce the effect of noise during analog-to-digital conversion is the Schmitt trigger²⁰, whereby two different analog thresholds are defined for the rising and falling edges of the discretized signal. For the case of a periodic input signal and colored Gaussian noise, the ocurrance of SR in this system has been demonstrated analytically²¹.

In dithering²², noise is added to audio or image signals in order to reduce the perception of quantization errors and the appearance of spurious large-scale patterns that can result from analog-to-digital conversion. Such spurious large scale patterns are associated with positive autocorrelations in the digitized signal. Therefore, although the aim in dithering is not to increase information transmission, the subjective perception of the signal by a human recipient may be improved by adding noise that in this case reduces such positive autocorrelations.

Related to this is the phenomenon of ghost stochastic resonance, which emerges when nonlinear systems are driven by periodic signals with more than one frequency in the presence of noise. An example for this effect is the auditory perception of 'ghost' frequencies that are not present in the acoustic pressure wave. The uncovering of this phenomenon helped to understand the perception of pitch in complex sound signals and the so-called 'miss-ing fundamental illusion'^{23, 24}. Interestingly, when noise of increasing intensity is added to sound, the emergence of a ghost resonance coincides with a maximum of the output autocorrelation^{23, 24}. This, in turn is closely linked to the effect of coherence resonance²⁵, whereby a non-linear excitable system is driven by noise, resulting in a maximally autocorrelated system output at a certain optimum noise intensity.

Although SR can explain a number of adaptive processes in neural systems, in particular in auditory and visual signal processing²⁶, how exactly that may be implemented is a matter of debate. Even for known input signals, the use of the mutual information or related approaches seem daunting for a biological system, since calculating them requires mathematical operations that are hard to implement in neuronal networks. By contrast, the autocorrelation function can be easily implemented with neuronal networks using delay-lines and coincidence detectors, as proposed in ref. 27 and verified, for example, in the nucleus laminaris of birds (barn owls), where they serve to code interaural time differences²⁸. More recent work has shown that a surrogate of the autocorrelation function can even be computed by a single noisy neuron^{23, 24}. Therefore, adaptive SR based on output autocorrelation may be a major processing principle in neural sensory systems and perception, and may also be responsible for pathologic conditions including neuropathic pain or tinnitus²⁶.



Figure 4. Performance of the output autocorrelation method. Comparison of resonance curves using the mutual information (red) and the output autocorrelation (blue) for different continuous signals. (a) Projection of the 3-dimensional Roessler attractor in one dimension, asymmetric threshold. (b) Ornstein-Uhlenbeck process, symmetric threshold. (c) Wave files of recorded speech, music and natural sounds, symmetric threshold. The biologically important case of the leaky integrate-and-fire model is shown in (d) with sine wave with constant frequency and amplitude as input. The signal-to-threshold distance was 0.1 for all cases.

Methods

Mutual Information. The mutual information I(S; Y) quantifies the mutual dependence of two random variables *S* and *Y*.¹⁶. It determines how similar the joint distribution p(s, y) is compared to the product of the marginal distributions p(s)p(y)

$$I(S; Y) = \sum_{s = y} p(s, y) \log_2 \left(\frac{p(s, y)}{p(s)p(y)} \right), \tag{1}$$

For continuous random variables, the summation is replaced by a double integral:

$$I(S; Y) = \int_{S} \int_{Y} p(s, y) \log_2\left(\frac{p(s, y)}{p(s)p(y)}\right) dy \, ds,$$
(2)

where p(s, y) is now the joint probability density function of *S* and *Y*, and p(s) and p(y) are the marginal probability density functions. The natural unit of I(S; Y) is *bits*, however in some cases it is more convenient to divide the total mutual information by the time, or by the number of spikes within the observed spike train, and thus derive mutual information rates R(S; Y) measured in *bits* s^{-1} or *bits spike*⁻¹. The choice of the mutual information as an objective function for adaptive stochastic resonance is natural, because the fundamental purpose of any sensor is to transmit information into a subsequent information processing system. Indeed, it has been shown by several authors^{3, 12, 14} that, within the context of stochastic resonance, I(S; Y) as a function of the variance σ^2 of the added noise has a maximum that indicates the optimal level of noise.



Figure 5. Correlation between optimal noise level from mutual information versus output autocorrlation. Numerical data from 5 sensor models, 3 different input signal, and different signal-to-threshold distances. The optimal noise intensities found by maximizing the mutual information (x-axes) are plotted versus the optimal noise intensities found by maximizing the output autocorrelation (y-axes). Except for correlated bipolar strings in combination with a leaky integrate-and-fire sensor, the optimal noise levels from both methods are highly correlated, with an overall correlation coefficient of r = 0.960 and a mean squared error of e = 0.027.

Output autocorrelation. The output autocorrelation function of the time lag τ (more precisely, the autocorrelation coefficient²⁹) is defined as

$$C_{yy}(\tau) = \frac{\langle (y_{t+\tau} - \bar{y}) (y_t - \bar{y}) \rangle_t}{\langle (y_t - \bar{y})^2 \rangle_t},$$
(3)

where \overline{y} is the mean and $\langle \cdot \rangle_t$ indicates averaging over time. We note that for most applications (and discrete time steps) it is sufficient to consider only one time lag, i.e. $\tau = 1$. However, for more complex signals, e.g. streams of n-bit words, it might be beneficial to calculate $C_{yy}(\tau)$ for a number of different subsequent lag times. In order to derive a single objective value in the case of multiple time lags τ , the root mean square (rms²⁹) of the autocorrelation function

$$\operatorname{RMS}(C_{yy}) = \sqrt{\frac{1}{N_{\tau}} \sum_{\tau} (C_{yy}(\tau))^2}$$
(4)

is calculated, where N_{τ} is the total number of different time lages.

Input signals for numerical simulations. Both, synthetically generated as well as natural signals were used in numerical simulations. As an example of a discrete input signal, a correlated, bipolar string $s_t \in \{-1, +1\}$ was generated, in which the probability of successive values being identical was $\operatorname{Prob}(s_t = s_{t-1}) = 0.7$. As examples of continuous signals we used: first, a sine waveform signal with constant frequency and amplitude; second, an aperiodic time series derived from the variable x(t) of the Roessler attractor³⁰

$$\begin{aligned} \dot{x} &= -(y+z) \\ \dot{y} &= x+ay \\ \dot{z} &= b+(x-c)z, \end{aligned} \tag{5}$$

with parameters a = 0.15, b = 0.2 and c = 7.1; third, the aperiodic random Ornstein-Uhlenbeck process³¹ $\dot{x} = -\frac{1}{\tau}x + \epsilon \xi(t)$, where ξ is an independent normally distributed random variable, τ is the correlation time and ϵ is the noise amplitude; fourth, wave files of speech, music and natural sounds. All synthetically generated signals were computed by numerically integrating the differential equations using fourth order Runge Kutta method.

Sensor models for numerical simulations. Four different memory-less sensor models were implemented by combining symmetric and asymmetric thresholds, with discrete and continuous sensor output functions (Fig. 5). In the symmetric models, there exist two thresholds $+\theta$ and $-\theta$. Without added noise, the sensor output is zero for $|s_t| < \theta$, $s_t - \theta$ for $s_t > \theta$, and $s_t + \theta$ for $s_t < -\theta$. In the asymmetric models, there exists only a single positive threshold $\theta > 0$. Here, without added noise, the sensor output is zero for $s_t < \theta$ and $s_t - \theta$ for $s_t > = \theta$. We note that the analytical model described above belongs to the class of discrete symmetric models. As one example of the detectors with memory, were the output depends not only on the momentary input and added noise, but also on earlier internal states of the detector itself, we choose the leaky integrate-and-fire neuron model^{32, 33} with

$$\dot{x} = -\frac{1}{\tau_m}x + s$$

where *x* is the membrane potential, τ_m the mebrane time constant and s_t the input signal. If *x* crosses the threshold θ from below, an output spike is generated and the membrane potential is set to the resting potential x_r , which is chosen to be zero for simplicity.

References

- 1. Benzi, R., Sutera, A. & Vulpiani, A. The mechanism of stochastic resonance. *Journal of Physics A: mathematical and general* 14, L453 (1981).
- Collins, J., Imhoff, T. & Grigg, P. Noise-enhanced information transmission in rat sa1 cutaneous mechanoreceptors via aperiodic stochastic resonance. *Journal of Neurophysiology* 76, 642–645 (1996).
- Levin, J. E. & Miller, J. P. Broadband neural encoding in the cricket cereal sensory system enhanced by stochastic resonance. *Nature* 380, 165–168 (1996).
- 4. Gammaitoni, L., Hänggi, P., Jung, P. & Marchesoni, F. Stochastic resonance. Reviews of modern physics 70, 223 (1998).
- 5. Wiesenfeld, K. & Jaramillo, F. Minireview of stochastic resonance. *Chaos: An Interdisciplinary Journal of Nonlinear Science* 8, 539–548 (1998).
- 6. Spezia, S. *et al.* Evidence of stochastic resonance in the mating behavior of nezara viridula (l.). *The European Physical Journal B* **65**, 453–458 (2008).
- 7. Hänggi, P. Stochastic resonance in biology how noise can enhance detection of weak signals and help improve biological information processing. *ChemPhysChem* **3**, 285–290 (2002).
- 8. Faisal, A. A., Selen, L. P. & Wolpert, D. M. Noise in the nervous system. Nature Reviews Neuroscience 9, 292-303 (2008).
- 9. Mino, H. The effects of spontaneous random activity on information transmission in an auditory brain stem neuron model. *Entropy* **16**, 6654–6666 (2014).
- Douglass, J. K., Wilkens, L., Pantazelou, E. & Moss, F. et al. Noise enhancement of information transfer in crayfish mechanoreceptors by stochastic resonance. Nature 365, 337–340 (1993).
- 11. Wiesenfeld, K. & Moss, F. *et al.* Stochastic resonance and the benefits of noise: from ice ages to crayfish and squids. *Nature* **373**, 33–36 (1995).
- Moss, F., Ward, L. M. & Sannita, W. G. Stochastic resonance and sensory information processing: a tutorial and review of application. *Clinical neurophysiology* 115, 267–281 (2004).
- 13. Mitaim, S. & Kosko, B. Adaptive stochastic resonance. *Proceedings of the IEEE* **86**, 2152–2183 (1998).
- Mitaim, S. & Kosko, B. Adaptive stochastic resonance in noisy neurons based on mutual information. *IEEE Transactions on Neural Networks* 15, 1526–1540 (2004).
- 15. Wenning, G. & Obermayer, K. Activity driven adaptive stochastic resonance. Physical review letters 90, 120602 (2003).
- 16. MacKay, D. J. Information theory, inference and learning algorithms (Cambridge university press, 2003).
- 17. Collins, J., Chow, C. C. & Imhoff, T. T. et al. Stochastic resonance without tuning. Nature 376, 236-238 (1995).
- 18. Greenwood, P. E., Müller, U. U. & Ward, L. M. Soft threshold stochastic resonance. Physical Review E 70, 051110 (2004).
- 19. Müller, U. U. & Ward, L. M. Stochastic resonance in a statistical model of a time-integrating detector. *Physical Review E* 61, 4286 (2000).
- 20. Schmitt, O. H. A thermionic trigger. Journal of Scientific Instruments 15, 24 (1938).
- 21. Melnikov, V. I. Schmitt trigger: A solvable model of stochastic resonance. Phys. Rev. E 48, 2481-2489 (1993).
- 22. Ken Pohlman, C. Principles of digital audio (2005).
- Chialvo, D. R. How we hear what is not there: A neural mechanism for the missing fundamental illusion. Chaos: An Interdisciplinary Journal of Nonlinear Science 13, 1226–1230 (2003).
- 24. Balenzuela, P., Braun, H. & Chialvo, D. R. The ghost of stochastic resonance: an introductory review. *Contemporary Physics* 53, 17–38 (2012).
- 25. Pikovsky, A. S. & Kurths, J. Coherence resonance in a noise-driven excitable system. Physical Review Letters 78, 775 (1997).
- 26. Krauss, P. *et al.* Stochastic resonance controlled upregulation of internal noise after hearing loss as a putative cause of tinnitus-related neuronal hyperactivity. *Frontiers in Neuroscience* **10** (2016).
- 27. Licklider, J. A duplex theory of pitch perception. The Journal of the Acoustical Society of America 23, 147-147 (1951).
- Sullivan, W. & Konishi, M. Neural map of interaural phase difference in the owl's brainstem. Proceedings of the National Academy of Sciences 83, 8400–8404 (1986).
- 29. Patrick, F. Measurement and data analysis for engineering and science (2005).
- 30. Rössler, O. E. An equation for continuous chaos. Physics Letters A 57, 397-398 (1976).
- 31. Ornstein, L. S. On the theory of the brownian motion. Physical Review 36, 823-841 (1930).
- Lapique, L. Recherches quantitatives sur l'excitation electrique des nerfs traitee comme une polarization. J Physiol Pathol Gen 9, 620–635 (1907).
- Burkitt, A. N. A review of the integrate-and-fire neuron model: I. homogeneous synaptic input. *Biological cybernetics* 95, 1–19 (2006).

Acknowledgements

We acknowledge support by Deutsche Forschungsgemeinschaft and Friedrich-Alexander-Universität Erlangen-Nürnberg (FAU) within the funding programme Open Access Publishing.

Author Contributions

P.K. and C.M. performed the calculations and computer simulations. P.K., C.M., K.T. and H.S. discussed the theoretical background. P.K., A.S. and C.S. prepared the figures. P.K., K.T., C.M., B.F. and H.S. wrote the manuscript.

Additional Information

Supplementary information accompanies this paper at doi:10.1038/s41598-017-02644-w

Competing Interests: The authors declare that they have no competing interests.

Publisher's note: Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.

Open Access This article is licensed under a Creative Commons Attribution 4.0 International License, which permits use, sharing, adaptation, distribution and reproduction in any medium or format, as long as you give appropriate credit to the original author(s) and the source, provide a link to the Creative Commons license, and indicate if changes were made. The images or other third party material in this article are included in the article's Creative Commons license, unless indicated otherwise in a credit line to the material. If material is not included in the article's Creative Commons license and your intended use is not permitted by statutory regulation or exceeds the permitted use, you will need to obtain permission directly from the copyright holder. To view a copy of this license, visit http://creativecommons.org/licenses/by/4.0/.

© The Author(s) 2017

Adaptive stochastic resonance for unknown and variable input signals

Patrick Krauss^{1,2}, Claus Metzner², Achim Schilling^{1,2}, Christian Schütz¹, Konstantin Tziridis¹, Ben Fabry² and Holger Schulze^{1,3}

¹Experimental Otolaryngology, ENT-Hospital, Head and Neck Surgery, Friedrich-Alexander University Erlangen-Nürnberg (FAU), Germany
²Department of Physics, Center for Medical Physics and Technology, Biophysics Group, Friedrich-Alexander University Erlangen-Nürnberg (FAU), Germany
³Correspondence to holger.schulze@uk-erlangen.de

April 5, 2017

Supplements

Success probability The output of a memory-less sensor can be described by a conditional probability distribution $p(y_t|s_t, n_t)$, which includes deterministic behaviour as a special case. Assuming statistically independent noise with distribution $p(n_t)$, the signal transmission properties of the sensor are given by $p(y_t|s_t) = \sum_{n_t} p(y_t|s_t, n_t)p(n_t)$. Ideally, the sensor output should be equal to the input signal, $y_t = s_t$, so that $p(y_t|s_t) = \delta_{y_t,s_t}$. It is therefore meaningful to quantify the performance of a sensor by the success probability

$$Q = p(y_t = s_t),\tag{1}$$

which is expected to peak at some optimum noise level within the context of stochastic resonance.

Analytical model In general, the momentary response of a SR-sensor can depend on the history of internal states of the system, as is the case in integrateand-fire-neurons. For simplicity, in the analytical model we only consider memoryless sensors, which respond to the present input signal s_t and noise value n_t independently from their former activity states.

We consider a bipolar stochastic sensor in which both the input signal s_t and the sensor output y_t can only take on the values -1 and +1. The noise values n_t , however, are continuous gaussian random numbers with variance σ^2 and mean $\mu = 0$, without any temporal correlations. We further assume that these two values appear in the input signal with equal probability, $p(s_t = -1) =$ $p(s_t = +1) = 0.5$. By assuming two symmetric detection thresholds $\pm \theta$ (figure 3 lower right inset), together with symmetric white noise, it can be assured that the distribution of sensor outputs $p(y_t = -1) = p(y_t = +1) = 0.5$ is also symmetric, so that the mean, variance and entropy of y_t remain constant even if the noise level is changed. Hence, the expressions for I(S; Y) and $C_{yy}(\tau)$ can be slightly simplified. In particular, the autocorrelation can be reduced to the non-normalized form $C_{yy}(\tau) \propto \langle y_t y_{t+\tau} \rangle$, and, furthermore, will be considered only for lagtime $\tau = 1$.

The sensor adds the noise n_t to the binary input signal s_t . If s_t+n_t exceeds the upper threshold θ , the output y_t is +1, if s_t+n_t falls below the lower threshold $-\theta$, output y_t is -1. For $s_t+n_t \in [-\theta, +\theta]$, the output is chosen randomly between the two binary values +1 and -1.

We are interested in the case of a threshold $\theta > 1$ which exceeds the signal amplitude, so that without the assistance of added noise the signal cannot be detected. Adding a random noise value n_t to a (say) positive input signal s_t can have three possible effects. If we consider the noise to be sufficiently positive to lift the signal beyond the upper threshold, then the success probability Q = $p(y_t = +1 | s_t = +1) = p(y_t = -1 | s_t = -1)$ will be increased. Alternatively, if the noise happens to be strongly negative and draws the positive signal below the lower threshold $-\theta$ then the success probability Q will be decreased. The third possibility is that s_t+n_t remains sub-threshold. Such cases make the signal transmission neither better nor worse. It is intuitively clear that small noise levels will increase Q, but as soon as a considerable fraction of momentary noise levels n_t exceeds $2 + (\theta - 1)$, the success probability Q will fall again. In our case it is given by $Q = Q(\sigma) = \frac{1}{2} + \frac{1}{2} \left[W(\frac{\theta+1}{\sigma}) - W(\frac{\theta-1}{\sigma}) \right]$, where $W(x) = \frac{1}{2} \operatorname{erf}(\frac{x}{\sqrt{2}})$ is a slightly rescaled error function (see *Derivation of success probability* for a detailed derivation). As a function of the noise level σ , the success probability has a well-defined maximum.

In this sensor model, the mutual information I(Y; S) can be expressed as a strictly increasing function of the success probability: $I(Q) = 1 + Q \log_2 Q + (1 - Q) \log_2(1-Q)$ (see *Derivation of mutual information* for a detailed derivation).

Since both I and Q require access to the sub-threshold signal s_t , we turn to the autocorrelation function C_{yy} of the sensor output. Since the mean \overline{y} of y_t is zero and its variance constant, we can use a non-normalized version of equation(6). Furthermore, we restrict our analytical consideration to a single lag-time $\tau = 1$, defining $C = \langle y_t y_{t+1} \rangle$. The modulus of this quantity, too, can be expressed as a strictly increasing function of the success probability: $|C(Q)| = |\langle s_t s_{t+1} \rangle| [1 - 4Q(1-Q)]$, where $\langle s_t s_{t+1} \rangle$ are the input correlations (see Derivation of output autocorrelation for a detailed derivation).

Derivation of success probability The normalized Gaussian distribution with zero-mean and standard deviation σ is given by

$$g(x,\sigma) = \frac{1}{\sqrt{2\pi\sigma}} e^{-\frac{1}{2}(x/\sigma)^2}$$
(2)

For later convenience, we define a function W(x) via

$$W\left(\frac{z}{\sigma}\right) = \int_0^z g(x,\sigma) \, dx = \frac{1}{2} \operatorname{erf}\left(\frac{1}{\sqrt{2}} \frac{z}{\sigma}\right),\tag{3}$$

where $\operatorname{erf}(x) = \frac{2}{\sqrt{\pi}} \int_0^x e^{-t^2} dt$ is the error function.

The success probability Q is given by

$$Q = p(y_t = +1|s_t = +1) = = \frac{1}{2} \cdot p(-\theta - 1 < n_t < \theta - 1) + + p(n_t > \theta - 1)$$
(4)

The factor $\frac{1}{2}$ accounts for the stochastic output of the unit in the case when $s_t + n_t$ is sub-threshold. We can now express the probabilities as integrals over Gaussians:

$$Q = \frac{1}{2} \cdot \left(\int_0^{\theta-1} g(x,\sigma) dx + \int_0^{\theta+1} g(x,\sigma) dx \right) + \left(\frac{1}{2} - \int_0^{\theta-1} g(x,\sigma) dx \right)$$
(5)

Next we use the function W(x) defined above:

$$Q = \frac{1}{2} \cdot \left(W\left(\frac{\theta-1}{\sigma}\right) + W\left(\frac{\theta+1}{\sigma}\right) \right) + \left(\frac{1}{2} - W\left(\frac{\theta-1}{\sigma}\right) \right) = \frac{1}{2} + \left[W\left(\frac{\theta+1}{\sigma}\right) - W\left(\frac{\theta-1}{\sigma}\right) \right]$$
(6)

Derivation of mutual information The mutual information of the detector output and the input signal is defined as

$$I(Y;S) = \sum_{y,s} p(y,s) \log_2\left(\frac{p(y,s)}{p(y)p(s)}\right) =$$

=
$$\sum_{y,s} p(y|s)p(s) \log_2\left(\frac{p(y|s)p(s)}{p(y)p(s)}\right) =$$

=
$$\sum_{y,s} p(y|s)(1/2) \log_2\left(\frac{p(y|s)(1/2)}{(1/2)(1/2)}\right) =$$

=
$$\frac{1}{2} \sum_{y,s} p(y|s) \log_2(2p(y|s)).$$
(7)

We explicitly go through all four terms:

$$2I(Y;S) = \sum_{y,s} p(y|s) \log_2 (2p(y|s)) =$$

$$= p(y=-1|s=-1) \log_2 (2p(y=-1|s=-1)) +$$

$$+ p(y=-1|s=+1) \log_2 (2p(y=-1|s=+1)) +$$

$$+ p(y=+1|s=-1) \log_2 (2p(y=+1|s=-1)) +$$

$$+ p(y=+1|s=+1) \log_2 (2p(y=+1|s=+1)) =$$

$$= Q \log_2 (2Q) +$$

$$+ (1-Q) \log_2 (2(1-Q)) +$$

$$+ (1-Q) \log_2 (2(1-Q)) +$$

$$+ Q \log_2 (2Q).$$
(8)

Therefore

$$I(Y;S) = Q \log_2(2Q) + (1-Q) \log_2(2(1-Q))$$

= 1+Q log₂(Q) + (1-Q) log₂(1-Q). (9)

Derivation of output autocorrelations in the analytical model The temporal correlations of the input signal can be expressed by the probability $q = p(s_1 = +1, s_0 = +1)$ in the following way:

$$\langle s_{t+1} \ s_t \rangle = \langle s_1 \ s_0 \rangle =$$

$$= \sum_{s_0, s_1} p(s_1, s_0) \ (s_1 \ s_0) =$$

$$= p(s_1 = -1 | s_0 = -1) p(s_0 = -1) \ [(-1)(-1)] +$$

$$+ p(s_1 = -1 | s_0 = +1) p(s_0 = +1) \ [(-1)(+1)] +$$

$$+ p(s_1 = +1 | s_0 = -1) p(s_0 = -1) \ [(+1)(-1)] +$$

$$+ p(s_1 = +1 | s_0 = +1) p(s_0 = +1) \ [(+1)(+1)] =$$

$$= q \ (1/2) \ [1] +$$

$$+ \ (1 - q) \ (1/2) \ [-1] +$$

$$+ \ q \ (1/2) \ [1] = 2q - 1.$$

$$(10)$$

The temporal correlations in the output signal are given by

$$C_{yy}(\tau = 1) = \langle y_{t+1} \ y_t \rangle = \langle y_1 \ y_0 \rangle =$$

=
$$\sum_{y_0, y_1} p(y_1, y_0) \ (y_1 \ y_0).$$
(11)

Consider for example the probability $p(y_1 = +1, y_0 = +1)$. There are four different chains of events which can produce a sequence of two successive +1's in the output signal:

$$p(y_{1} = +1, y_{0} = +1) =$$

$$= p(y_{1} = +1|s_{1} = -1)p(s_{1} = -1|s_{0} = -1) \cdot$$

$$\cdot p(y_{0} = +1|s_{0} = -1)p(s_{0} = -1) +$$

$$+ p(y_{1} = +1|s_{1} = -1)p(s_{1} = -1|s_{0} = +1) \cdot$$

$$\cdot p(y_{0} = +1|s_{0} = +1)p(s_{0} = +1) +$$

$$+ p(y_{1} = +1|s_{1} = +1)p(s_{1} = +1|s_{0} = -1) \cdot$$

$$\cdot p(y_{0} = +1|s_{0} = -1)p(s_{0} = -1) +$$

$$+ p(y_{1} = +1|s_{1} = +1)p(s_{1} = +1|s_{0} = +1) \cdot$$

$$\cdot p(y_{0} = +1|s_{0} = +1)p(s_{0} = +1) =$$

$$= (1 - Q) q (1 - Q) (1/2) +$$

$$+ (1 - Q) (1 - Q) (1/2) +$$

$$+ Q (1 - q) (1 - Q) (1/2) +$$

$$+ Q q Q (1/2) =$$

$$= \frac{q}{2} + (2q - 1)Q(1 - Q) =: A.$$
(12)

For symmetry reasons, $p(y_1 = -1, y_0 = -1) = p(y_1 = +1, y_0 = +1) = A$. In the same way, $p(y_1 = +1, y_0 = -1) = p(y_1 = -1, y_0 = +1) = B$.

Since
$$\sum_{y_0, y_1} p(y_1, y_0) = 1 = 2A + 2B$$
, it follows that $B = \frac{1}{2} - A$.

Knowing all four joint probabilities, we can proceed to compute the temporal correlations in the output signal:

$$C_{yy}(\tau = 1) = \langle y_1 \ y_0 \rangle =$$

= $A(-1)(-1) + B(-1)(+1) +$
+ $B(+1)(-1) + A(+1)(+1) =$
= $2A - 2B =$
= $(2q - 1) [1 - 4Q(1 - Q)] =$
= $\langle s_{t+1} \ s_t \rangle [1 - 4Q(1 - Q)].$ (13)

Soft thresholds and non-Gaussian noise As described above, the probabilistic information transmission from the signal input s_t to the output y_t of a sensor are defined by

$$p(y_t|s_t) = \int_{-\infty}^{+\infty} p(y_t|s_t, n_t) \ p(n_t) \ dn_t.$$
(14)

Here, $p(y_t|s_t, n_t)$ characterizes the properties of a specific sensor type, and $p(n_t) = p_{noi}(n_t)$ is the noise distribution.

In this work, we are considering sensors where signal and noise are combined additively, so that $p(y_t|s_t, n_t)$ can be replaced by a simpler conditional probability that depends only on the sum $x_t = s_t + n_t$:

$$p(y_t|s_t, n_t) \longrightarrow p(y_t \mid s_t + n_t = x_t) = p(y_t|x_t).$$

$$(15)$$

Furthermore, in the case of bipolar sensors, where $y_t = +1$ and $y_t = -1$ are the only possible outputs, the sensor can be characterized by a *response function*:

$$P_{res}(x) = p(y_t = +1|x_t = x).$$
(16)

For such additive, bipolar sensors, the success probability Q can be expressed via the response function and the noise distribution as

$$Q = p(y_t = +1|s_t = +1)$$

= $\int_{-\infty}^{+\infty} P_{res}(x) p(n_t = x-1) dx$
= $\int_{-\infty}^{+\infty} P_{res}(x) p_{noi}(x-1) dx.$ (17)

So far in this Supplemental, we have only considered detectors with a piecewise constant respose function (that is, where $P_{res}(x < -\theta) = 0$ and $P_{res}(-\theta < = x < = +\theta) = \frac{1}{2}$ and $P_{res}(x > +\theta) = 1$), and where the noise was normal distributed (see Fig.1(a)).

However, we can easily generalize our analytical model to permit arbitrary response functions (for example, by using a smooth sigmoidal function rather than one with hard thresholds) and non-Gaussian noise distributions. In order to keep the symmetry $p(y_t = -1) = p(y_t = +1) = 0.5$ of the output signals (which we have used to simplify our analytical derivation), we have to restrict our choices to response functions with $P_{res}(-x) = 1 - P_{res}(+x)$ and to noise distributions with $p_{noi}(-x) = p_{noi}(+x)$. One of the possible choices is sketched in Fig.1(b).

The generalization of the model to smooth sigmoidal sensor responses and non-Gaussian noise does only affect the success probability $Q = Q(\sigma)$ and its dependence on the noise amplitude σ . As long as $Q(\sigma)$ has a peak at some optimum noise level σ_{opt} , the strictly monotonous dependence of the mutual information I and the output correlations C_{yy} on Q guaranty that I and C_{yy} will peak at the same noise level σ_{opt} .

Different types of detectors and multiplicative noise For all bipolar sensors (with $s_t \in \{-1, +1\}$ and $y_t \in \{-1, +1\}$) a success probability can always be defined as

$$Q = \int_{-\infty}^{+\infty} p(y_t = +1 \mid s_t = +1, n_t) \ p(n_t) \ dn_t.$$
(18)

The general conditional probability $p(y_t = +1 | s_t = +1, n_t)$ includes not only detectors where signal and noise are combined additively, but allows for an arbitrary probabilistic dependence of y_t on s_t and n_t . As long as the noise n_t is temporally un-correlated (and all variables s_t , n_t and y_t are zero-mean) C_{yy} and the MI will remain monotonous functions of Q.

However, it is not guaranteed that these objective functions will always have a maximum as a function of the noise strength σ . Consider, for example, a sensor system where the zero-mean, temporally correlated signal s_t is *multiplied* with zero-mean, white noise n_t , and where the product $x_t = s_t * n_t$ is compared with a hard or soft sigmoidal threshold, as above. In this case, the amplitude modulated noise x_t is also un-correlated, so that $C_{yy} = 0$, no matter how the noise strength σ is set. Thus, for such multiplicative systems, the output autocorrelation is not in general a suitable objective function for adaptive SR. In the context of neural systems, however, the assumptions of a symmetric bipolar threshold and of completely un-correlated noise are not biologically plausible.



Figure 1: Sketch of possible response functions (blue) and noise distributions (green). (a) Piecewise constant $P_{res}(x)$ and Gaussian $p_{noi}(x)$. (b) Smooth sigmoidal $P_{res}(x)$ and leptocurtic $p_{noi}(x)$